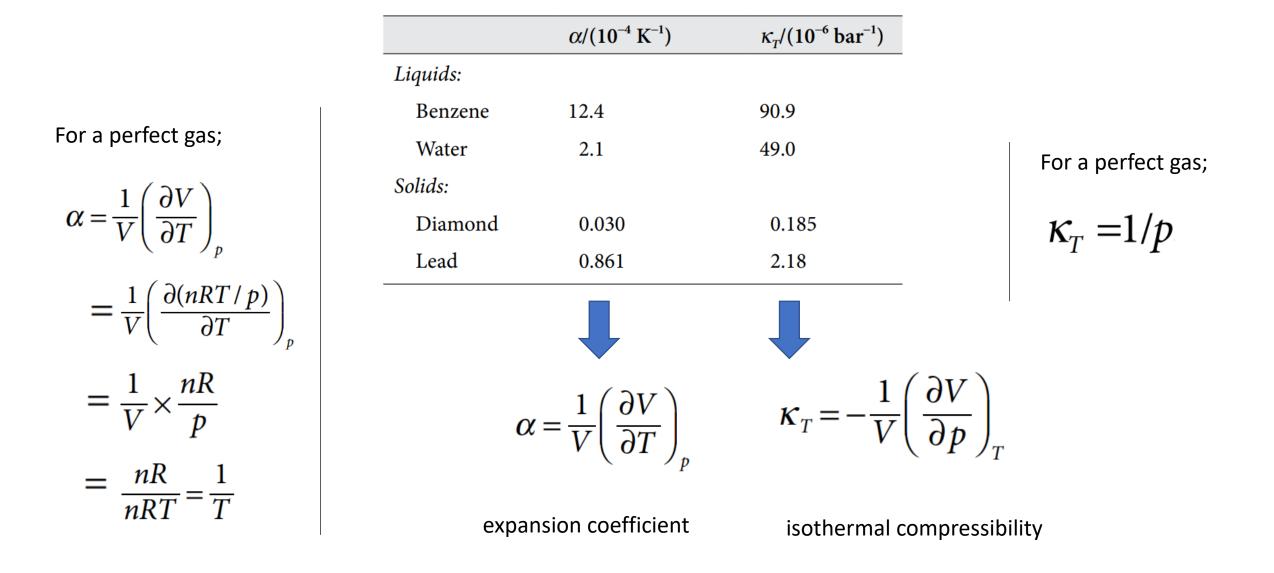
Expansion and isothermal coefficients





$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$

 $\mathrm{d}U = \pi_{T}\mathrm{d}V + C_{V}\mathrm{d}T$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

 $\mathrm{d}U = \pi_{T}\mathrm{d}V + C_{V}\mathrm{d}T$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

 $\mathrm{d}U = \pi_{T}\mathrm{d}V + C_{V}\mathrm{d}T$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

 $\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas, $\pi_T = 0$, $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

 $dU = \pi_T dV + C_V dT$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

for a perfect gas:

$$\left(\frac{dU}{dT}\right)_V = \left(\frac{dU}{dT}\right)_P = C_v$$

For a perfect gas, $\pi_T = 0$, $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

Since the internal energy U is a function of temperature only

$$\mathrm{d}U = \pi_{T}\mathrm{d}V + C_{V}\mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$
, $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

$$C_p - C_v =$$

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$
, $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

$$C_{p} - C_{v} = \underbrace{\left(\frac{\partial H}{\partial T}\right)_{p}}^{\text{Definition}} - \left(\frac{\partial U}{\partial T}\right)_{p}$$

$$\mathrm{d}U = \pi_{T}\mathrm{d}V + C_{V}\mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$

 $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

$$C_{p} - C_{v} = \underbrace{\left(\frac{\partial H}{\partial T}\right)_{p}}^{\text{Definition}} - \left(\frac{\partial U}{\partial T}\right)_{p}$$

H = U + pV

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$

 $\left(\frac{\partial U}{\partial T}\right)_p = C_V$

Definition
of
$$C_p$$

 $C_p - C_V = \left(\frac{\partial H}{\partial T} \right)_p - \left(\frac{\partial U}{\partial T} \right)_p$

H = U + pV = U + nRT

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$

$$C_{p} - C_{v} = \underbrace{\left(\frac{\partial H}{\partial T}\right)_{p}}^{\text{Definition}} - \left(\frac{\partial U}{\partial T}\right)_{p}$$

H = U + pV = U + nRT

$$C_p - C_V = \left(\frac{\partial (U + nRT)}{\partial T}\right)_p - \left(\frac{\partial U}{\partial T}\right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = C_V$$

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$

$$C_{p} - C_{v} = \left(\underbrace{\frac{\partial H}{\partial T}}_{p} \right)_{p} - \left(\frac{\partial U}{\partial T} \right)_{p}$$

H = U + pV = U + nRT

$$C_p - C_V = \left(\frac{\partial (U + nRT)}{\partial T}\right)_p - \left(\frac{\partial U}{\partial T}\right)_p = nR$$

$$\left(\frac{\partial U}{\partial T}\right)_p = C_V$$

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

For a perfect gas,
$$\pi_T = 0$$

$$C_{p} - C_{v} = \underbrace{\left(\frac{\partial H}{\partial T}\right)_{p}}^{\text{Definition}} - \left(\frac{\partial U}{\partial T}\right)_{p}$$

H = U + pV = U + nRT

For perfect gases

$$C_p - C_V = \left(\frac{\partial (U + nRT)}{\partial T}\right)_p - \left(\frac{\partial U}{\partial T}\right)_p = nR$$

$$\left(\frac{\partial U}{\partial T}\right)_p = C_V$$

$$\mathrm{d}U = \pi_T \mathrm{d}V + C_V \mathrm{d}T$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \pi_T \left(\frac{\partial V}{\partial T}\right)_p + C_V$$

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$

$$\left(\frac{\partial U}{\partial T}\right)_p = \alpha \pi_T V + C_V$$

$$= 0 \qquad \left(\frac{\partial U}{\partial T}\right)_p = C_V$$

For a perfect gas, $\pi_T = 0$

$$C_{p} - C_{v} = \underbrace{\left(\frac{\partial H}{\partial T}\right)_{p}}^{\text{Definition}} - \left(\frac{\partial U}{\partial T}\right)_{p}$$

H = U + pV = U + nRT

For perfect gases

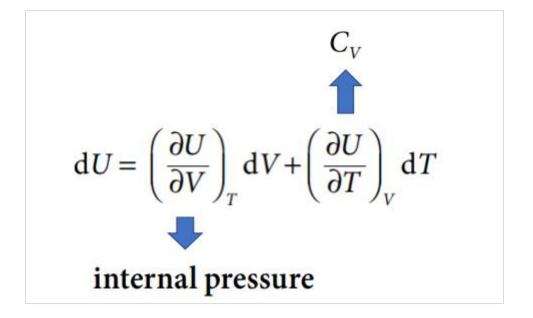
$$C_p - C_V = \left(\frac{\partial (U + nRT)}{\partial T}\right)_p - \left(\frac{\partial U}{\partial T}\right)_p = nR$$

$$C_p - C_V = \frac{\alpha^2 T V}{\kappa_T}$$

for any other gas

$$\mathbf{d}H = \left(\frac{\partial H}{\partial p}\right)_T \mathbf{d}p + \left(\frac{\partial H}{\partial T}\right)_p \mathbf{d}T$$

The variation of enthalpy with temperature and pressure



$$\mathrm{d}H = \left(\frac{\partial H}{\partial p}\right)_T \mathrm{d}p + \left(\frac{\partial H}{\partial T}\right)_p \mathrm{d}T$$

The variation of enthalpy with temperature and pressure

$$\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$$

$$\mathrm{d}H = \left(\frac{\partial H}{\partial p}\right)_T \mathrm{d}p + \left(\frac{\partial H}{\partial T}\right)_p \mathrm{d}T$$

The variation of enthalpy with temperature and pressure

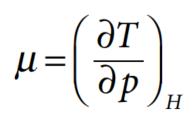
$$dH = -\mu C_p dp + C_p dT$$

Joule-Thomson coefficient

$$\mathrm{d}H = \left(\frac{\partial H}{\partial p}\right)_T \mathrm{d}p + \left(\frac{\partial H}{\partial T}\right)_p \mathrm{d}T$$

The variation of enthalpy with temperature and pressure

$$dH = -\mu C_p dp + C_p dT$$
Joule-Thomson coefficient



 $\mu = \left(\frac{\partial T}{\partial p}\right)_{H}$ the variation of temperature with pressure at constant enthalpy

Joule-Thomson effect

The Joule-Thomson effect (or Joule-Kelvin effect) describes a real gas's temperature change when it undergoes an isenthalpic process.

Joule-Thomson effect

The Joule-Thomson effect (or Joule-Kelvin effect) describes a real gas's temperature change when it undergoes an isenthalpic process.

LAW: Show that the Joule-Thomson (JT) effect is an isenthalpic process using gas flowing through a porous plug or a throttle.

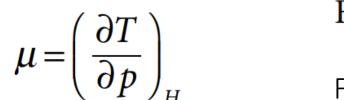
$$\mu = \left(\frac{\partial T}{\partial p}\right)_{H}$$

$$\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$$

For a perfect gas, $\mu = 0$

 $\mu = \left(\frac{\partial T}{\partial p}\right)_{H}$

 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$



For a perfect gas, $\mu = 0$

For, real gases, $\mu > 0$ or $\mu < 0$, depending on the conditions

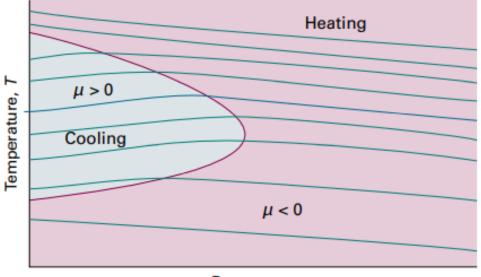
 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$

For a perfect gas, $\mu = 0$

For, real gases, $\mu > 0$ or $\mu < 0$, depending on the conditions

 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$

 $\mu = \left(\frac{\partial T}{\partial p}\right)_{\mu}$



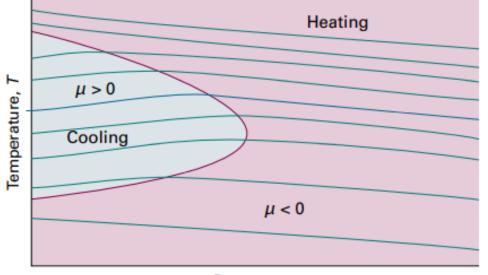
Pressure, p

For a perfect gas, $\mu = 0$

For, real gases, $\mu > 0$ or $\mu < 0$, depending on the conditions

 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$

 $\mu = \left(\frac{\partial T}{\partial p}\right)_{T}$



The temperature corresponding to the boundary at a given pressure is the 'inversion temperature' of the gas at that pressure

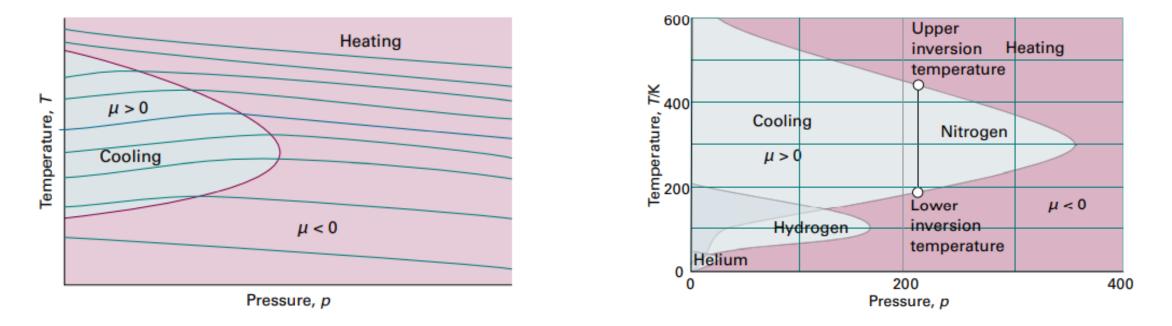
Pressure, p

For a perfect gas, $\mu = 0$

For, real gases, $\mu > 0$ or $\mu < 0$, depending on the conditions

 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$

 $\mu = \left(\frac{\partial T}{\partial p}\right)_{\mu}$



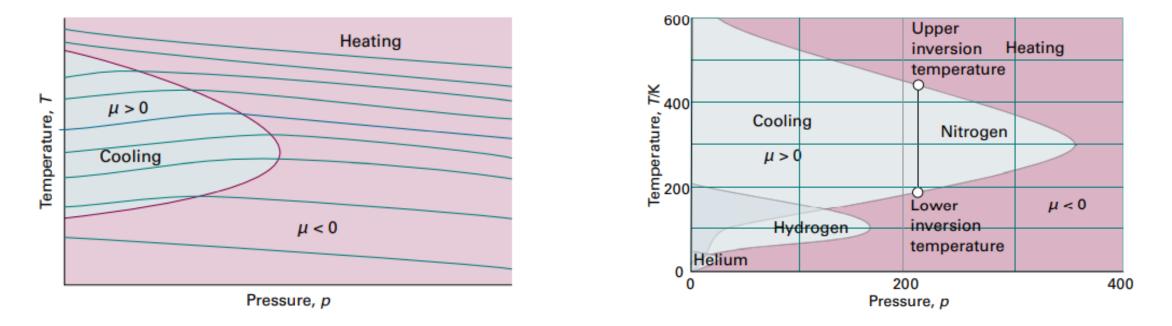
The inversion temperature curve runs through the points of the isenthalps where their slope changes from negative to positive.

For a perfect gas, $\mu = 0$

For, real gases, $\mu > 0$ or $\mu < 0$, depending on the conditions

 $\mathbf{d}H = -\mu C_p \mathbf{d}p + C_p \mathbf{d}T$

 $\mu = \left(\frac{\partial T}{\partial p}\right)_{\mu}$



The inversion temperature curve runs through the points of the isenthalps where their slope changes from negative to positive.

$$\mu = \left(\frac{\partial T}{\partial p}\right)_{H}$$

Nitrogen:

Has a high inversion temperature, allowing cooling over a wide temperature range.

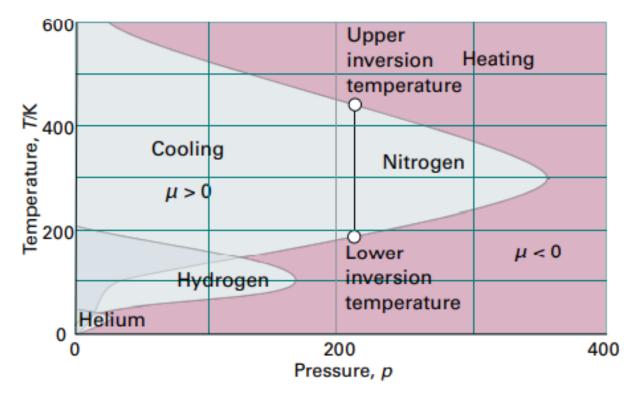
• Is easily liquefied using the Joule-Thomson effect.

Hydrogen and Helium:

•Both have very low inversion temperatures (below room temperature).

•At room temperature, expanding hydrogen or helium causes heating!

•To cool them via the JT effect, they must first be precooled below their inversion temperatures using other methods.



Focus 2: The First Law Internal Energy Enthalpy Thermochemistry State functions Adiabatic changes