

Temperature dependence of reaction enthalpies

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$$dH = C_p dT$$

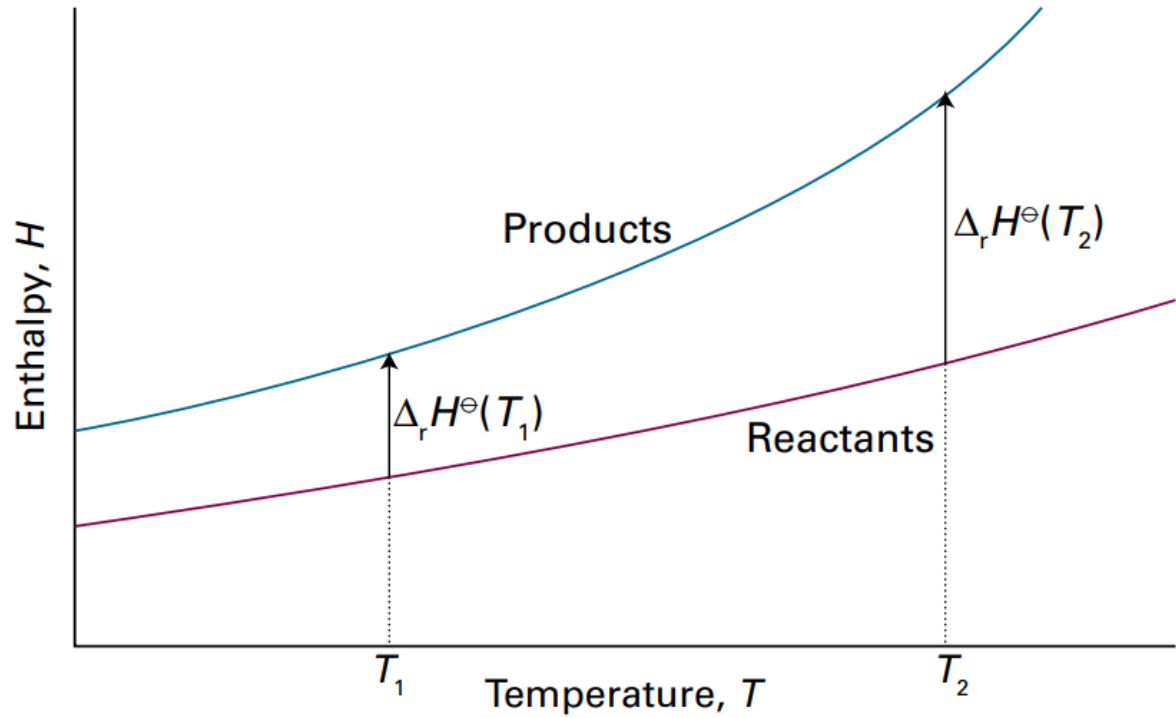
Temperature dependence of reaction enthalpies

$$dH = C_p dT$$

$$H(T_2) = H(T_1) + \int_{T_1}^{T_2} C_p dT$$

If no phase transition takes place

Temperature dependence of reaction enthalpies

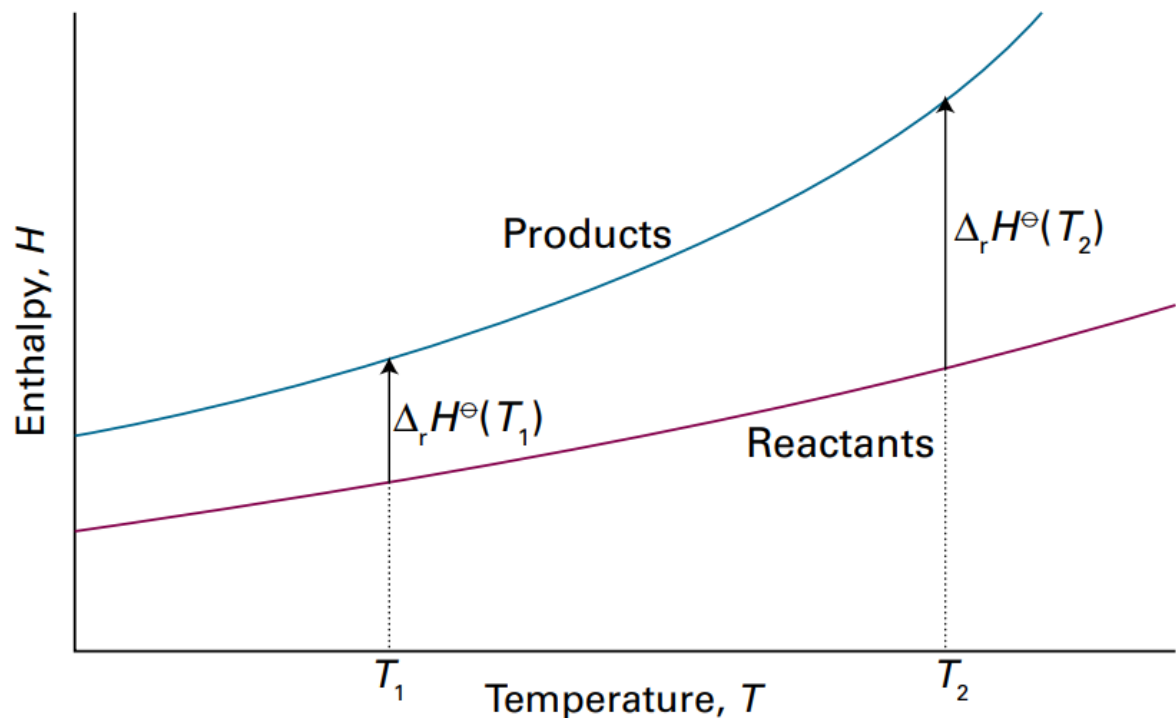


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Kirchhoff's law

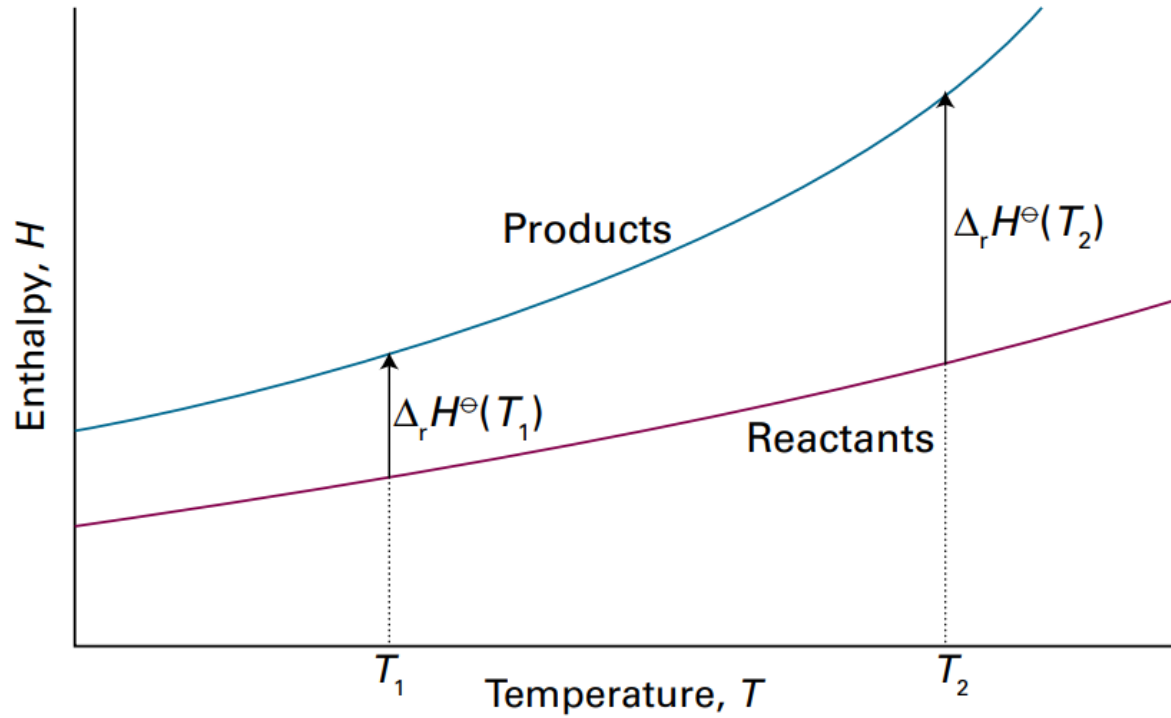
$$\Delta_r H^\ominus(T_2) = \Delta_r H^\ominus(T_1) + \int_{T_1}^{T_2} \Delta_r C_p^\ominus dT$$

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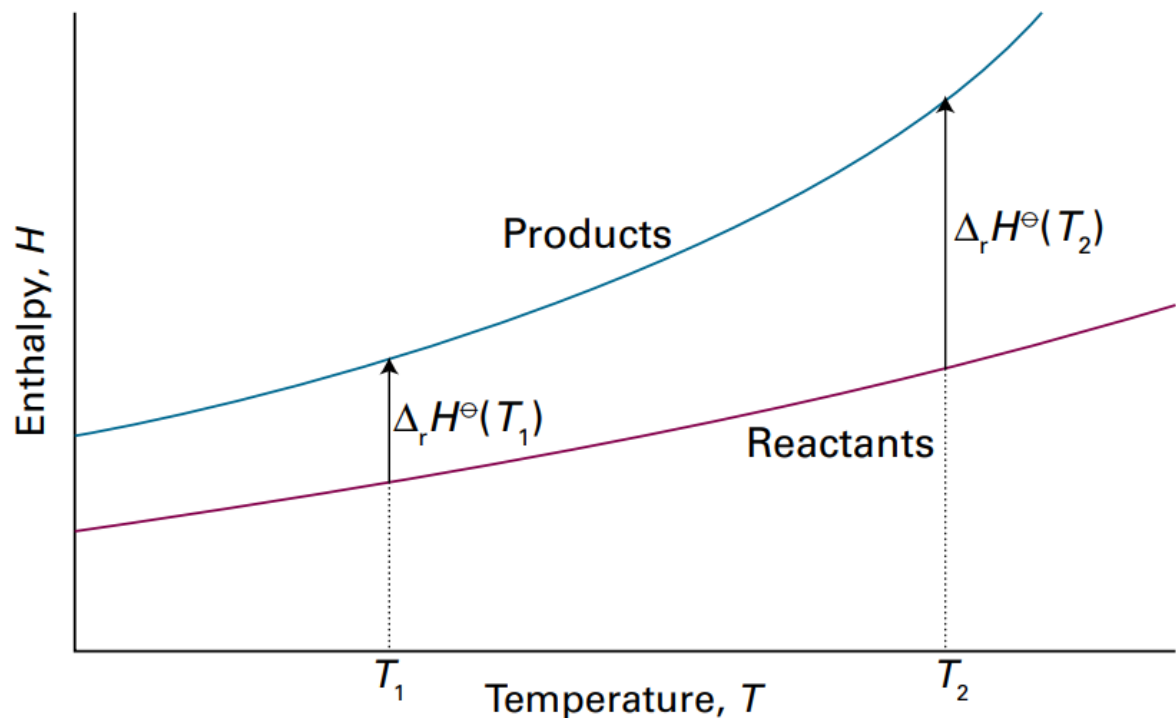
$$\Delta_r C_p^\ominus = \sum_{\text{Products}} \nu C_{p,m}^\ominus - \sum_{\text{Reactants}} \nu C_{p,m}^\ominus$$

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Integrated
form of
Kirchhoff's law

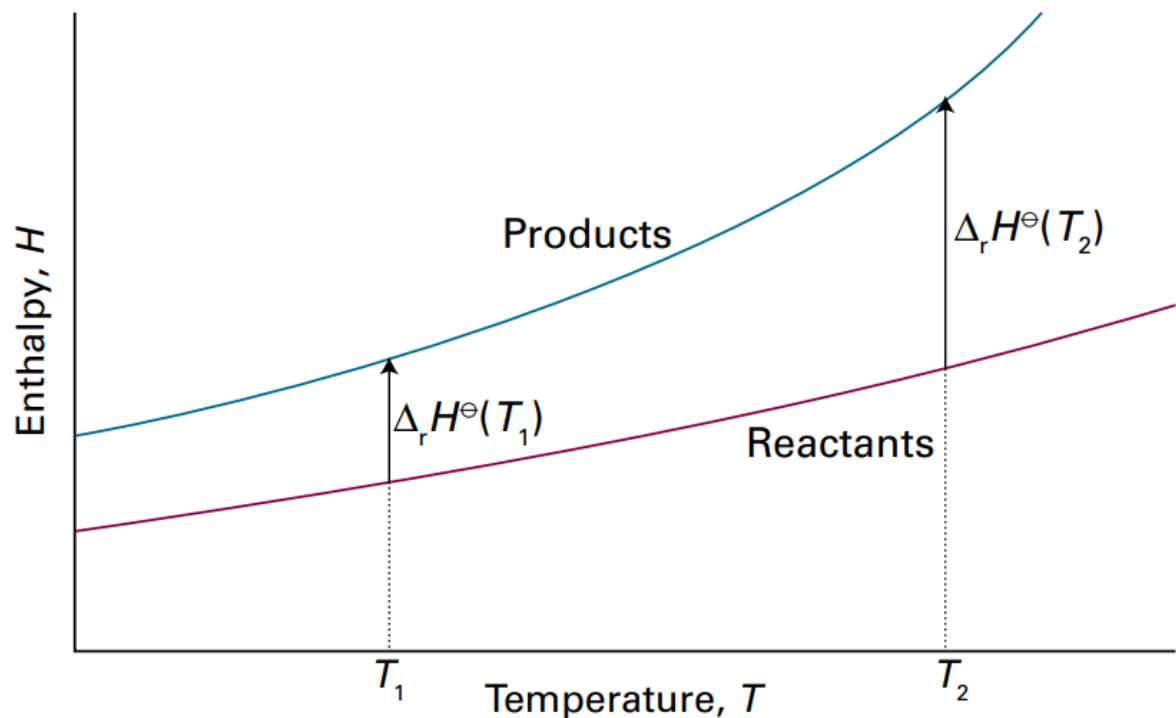
The difference in heat capacities
varies less significantly

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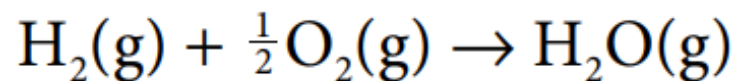
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Integrated form of Kirchhoff's law

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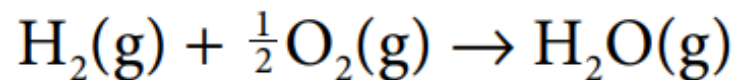
If no phase transition takes place

Calculation

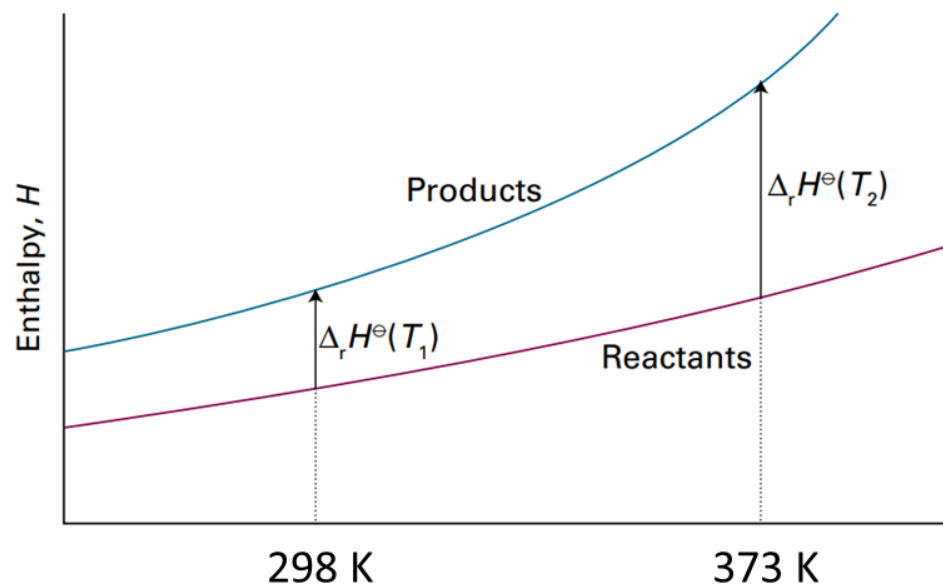


The standard enthalpy of formation of $\text{H}_2\text{O}(\text{g})$ at 298 K is $-241.82 \text{ kJ mol}^{-1}$. Estimate its value at 100°C given the following values of the molar heat capacities at constant pressure: $\text{H}_2\text{O}(\text{g})$: $33.58 \text{ J K}^{-1} \text{ mol}^{-1}$; $\text{H}_2(\text{g})$: $28.84 \text{ J K}^{-1} \text{ mol}^{-1}$; $\text{O}_2(\text{g})$: $29.37 \text{ J K}^{-1} \text{ mol}^{-1}$. Assume that the heat capacities are independent of temperature.

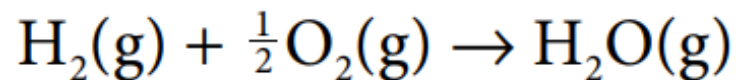
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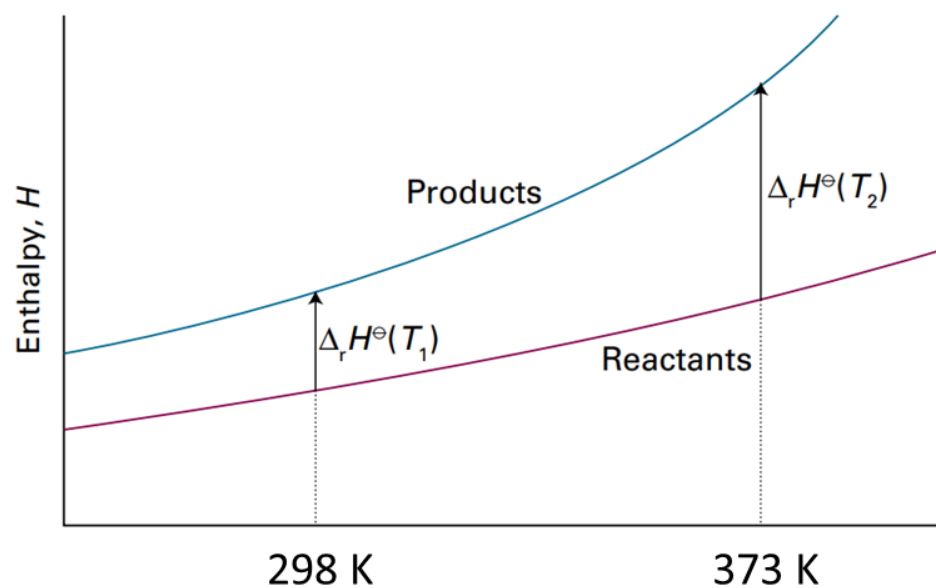
Calculation



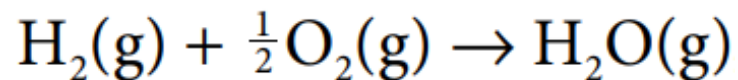
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Kirchhoff's law

$$\Delta_r H^\ominus(T_2) = \Delta_r H^\ominus(T_1) + \Delta_r C_p^\ominus(T_2 - T_1)$$



Calculation



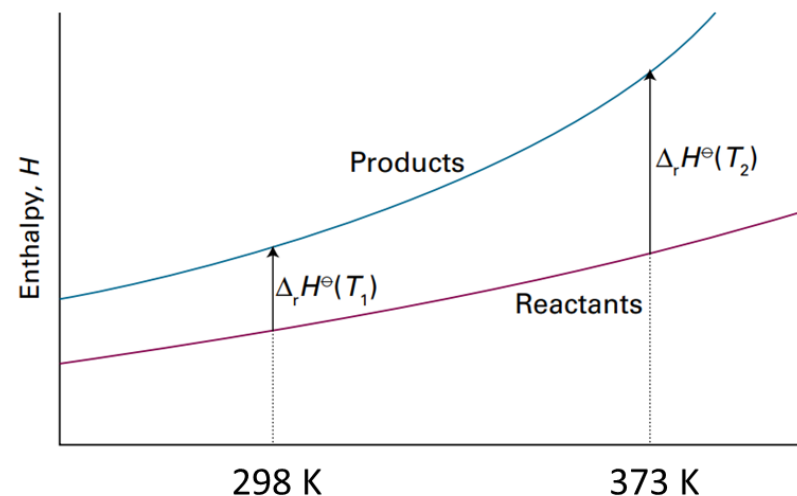
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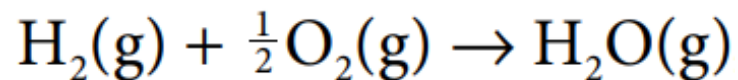
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↓

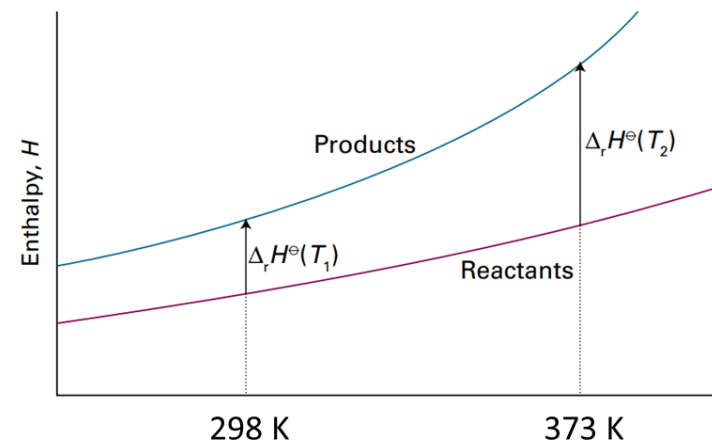
$$-241.82 \text{ kJ mol}^{-1}$$



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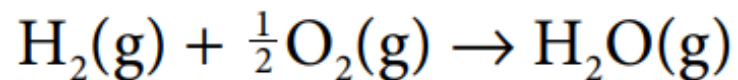
\downarrow

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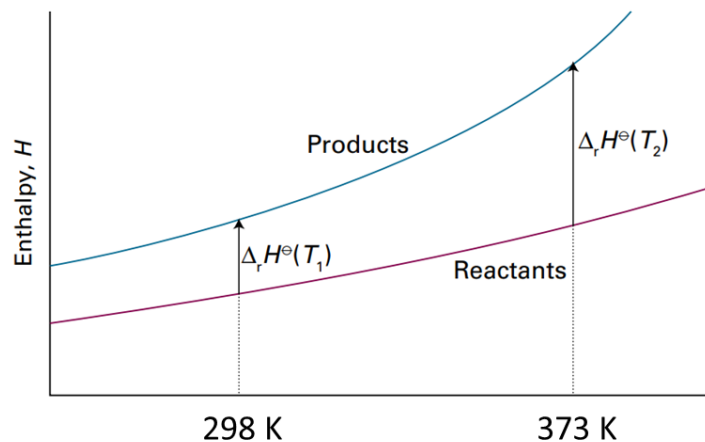
\downarrow

$$\begin{aligned}\Delta_r C_p^\ominus &= C_{p,m}^\ominus(\text{H}_2\text{O}, \text{g}) - \{C_{p,m}^\ominus(\text{H}_2, \text{g}) + \frac{1}{2} C_{p,m}^\ominus(\text{O}_2, \text{g})\} \\ &= -9.94 \text{ J K}^{-1} \text{ mol}^{-1}\end{aligned}$$

Calculation



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Kirchhoff's law

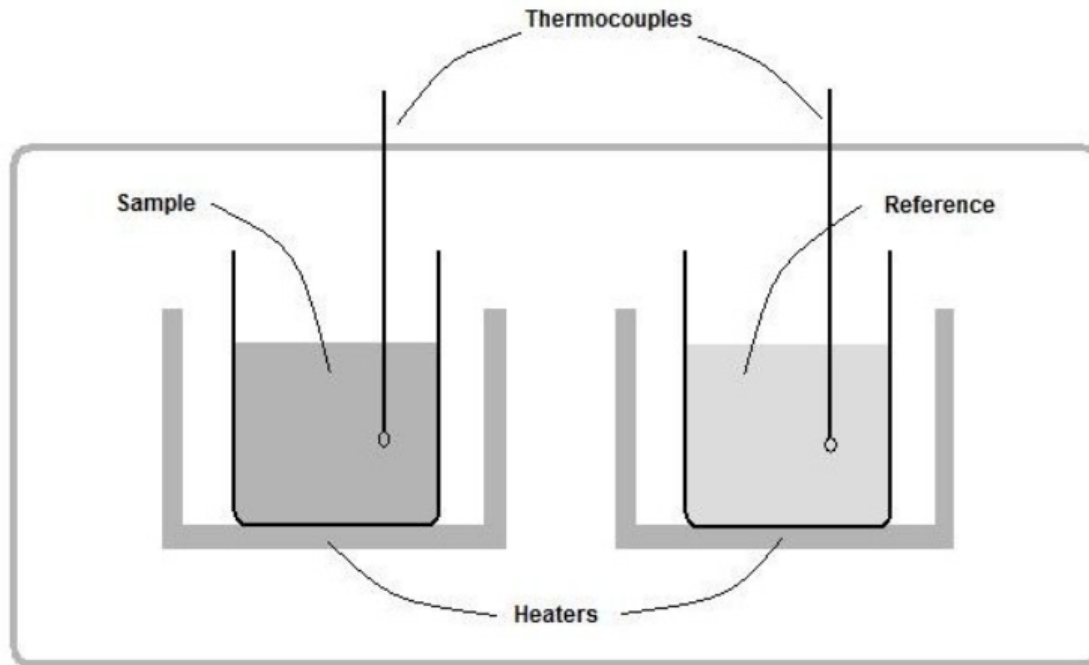
$$\Delta_r H^\ominus(T_2) = \Delta_r H^\ominus(T_1) + \Delta_r C_p^\ominus(T_2 - T_1)$$

\downarrow
 $-241.82 \text{ kJ mol}^{-1}$

\downarrow
$$\Delta_r C_p^\ominus = C_{p,m}^\ominus(\text{H}_2\text{O}, \text{g}) - \{C_{p,m}^\ominus(\text{H}_2, \text{g}) + \frac{1}{2} C_{p,m}^\ominus(\text{O}_2, \text{g})\}$$
$$= -9.94 \text{ J K}^{-1} \text{ mol}^{-1}$$

$$\Delta_r H^\ominus(373 \text{ K}) = -241.82 \text{ kJ mol}^{-1} + (75 \text{ K})$$
$$\times (-9.94 \text{ J K}^{-1} \text{ mol}^{-1}) = -242.6 \text{ kJ mol}^{-1}$$

Differential Scanning Calorimetry



- Sample and reference (solvent) reside in separate chambers
- Separate sources heat the chambers in a way that their temperatures are always equal
- The heating rates, which were used to maintain equivalent temperatures, are logged
- Any difference between the two can be attributed to the presence of the substance of interest

Focus 2: The First Law

Internal Energy

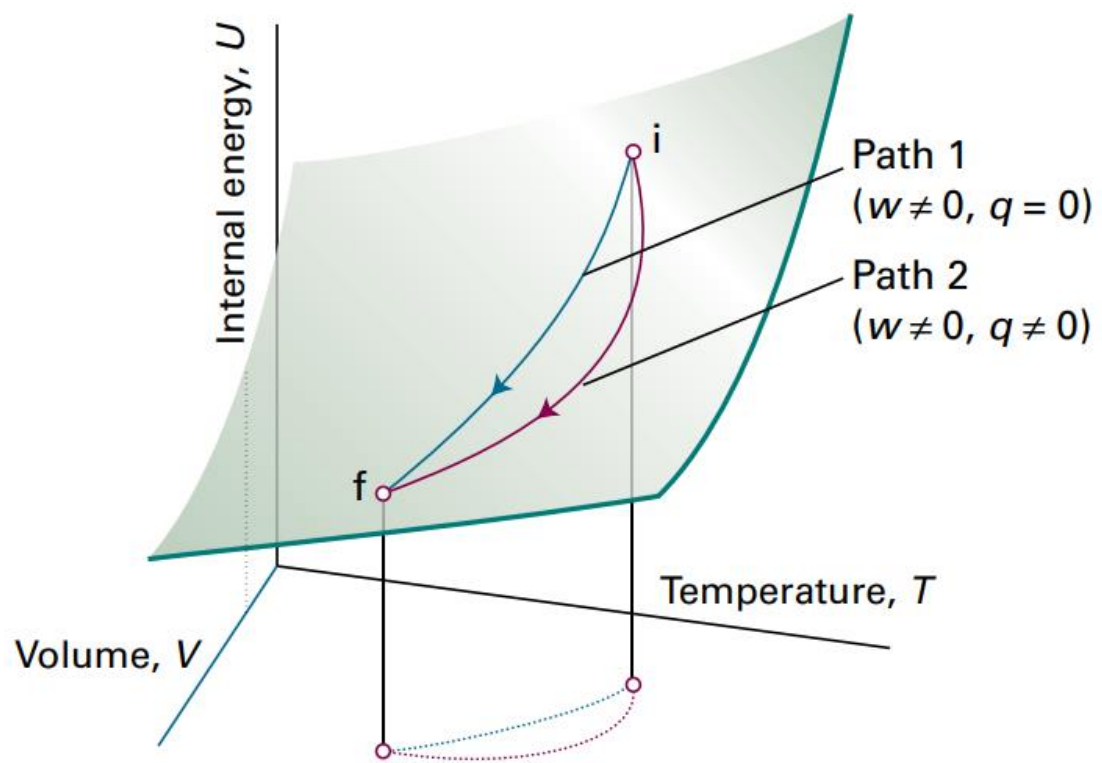
Enthalpy

Thermochemistry

State functions

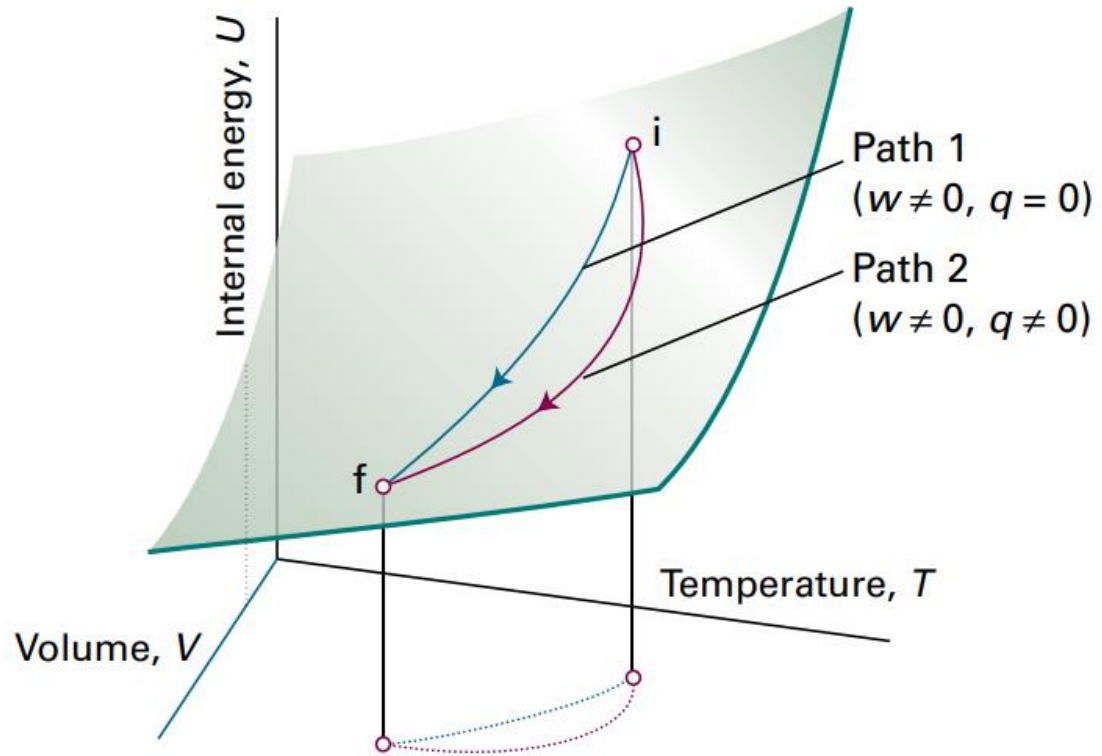
Adiabatic changes

Exact and inexact differentials



Exact and inexact differentials

$$\Delta U = \int_i^f dU$$



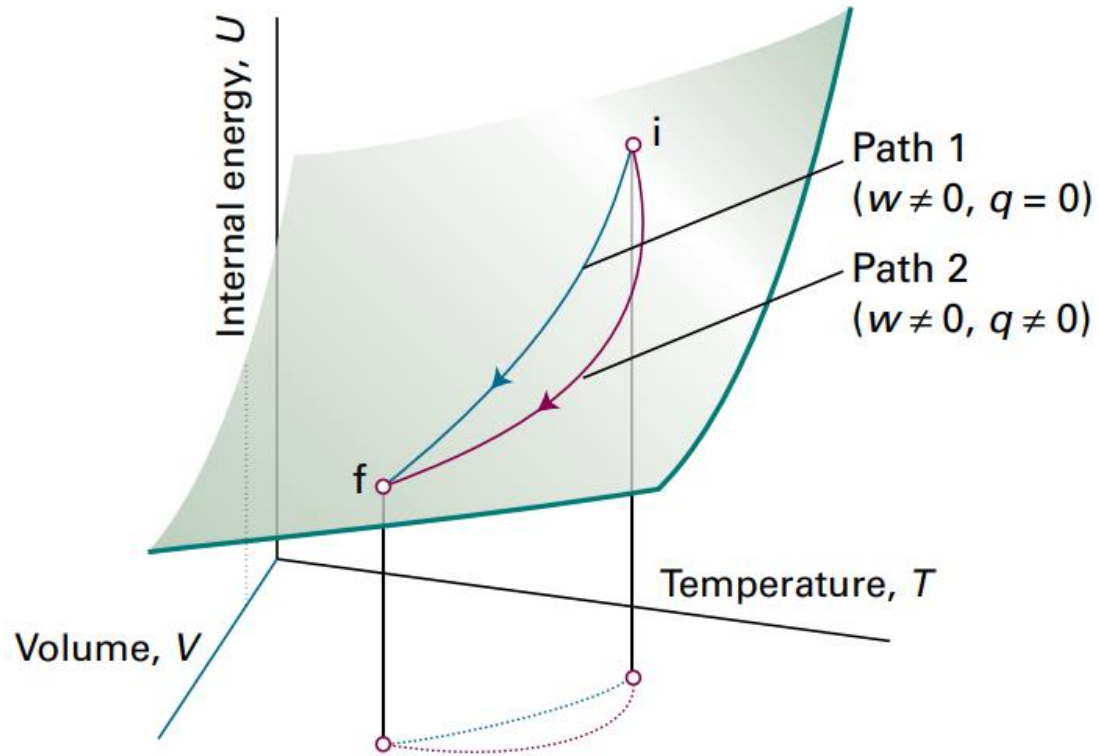
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exact differential

A differential dX is called exact if it represents the differential of a state function



Exact and inexact differentials

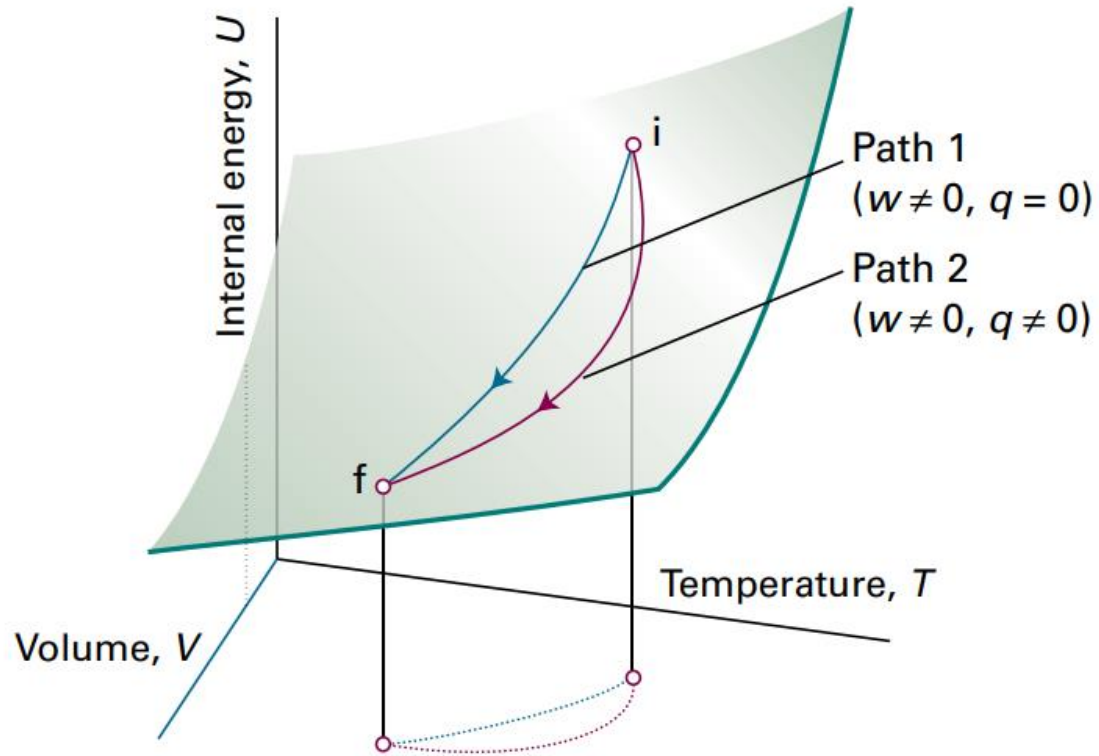
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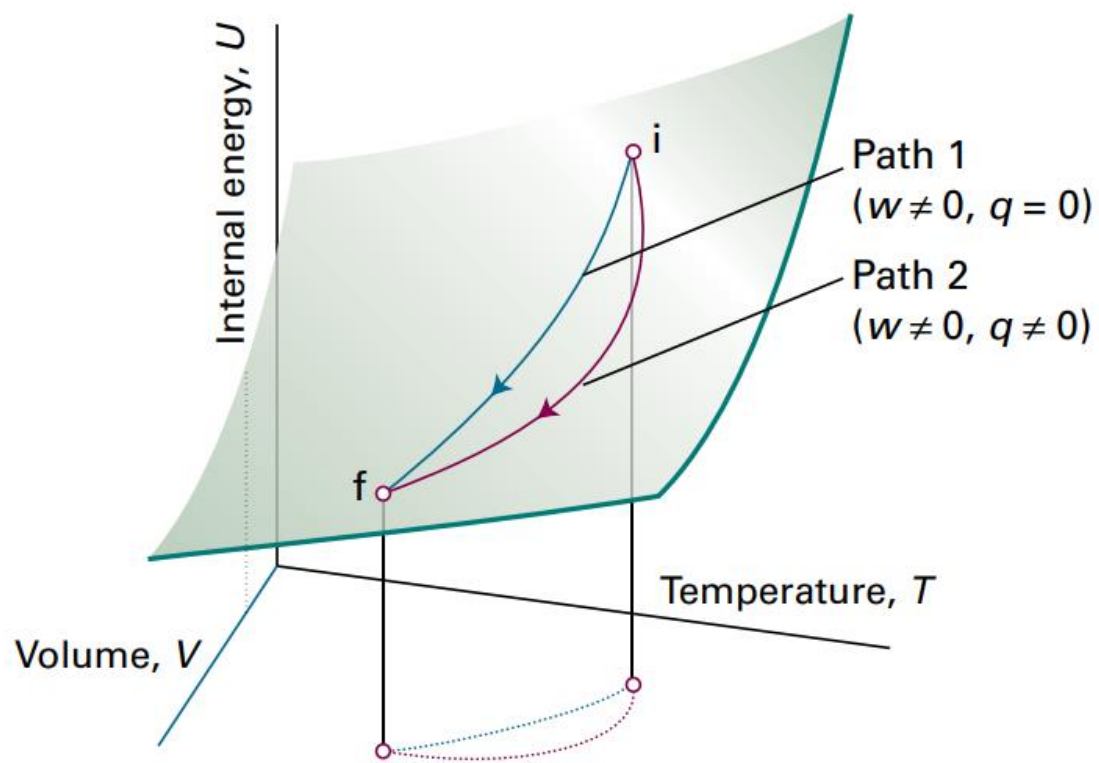
exact differential

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$$q = \int_{i, \text{path}}^f dq$$



Exact and inexact differentials



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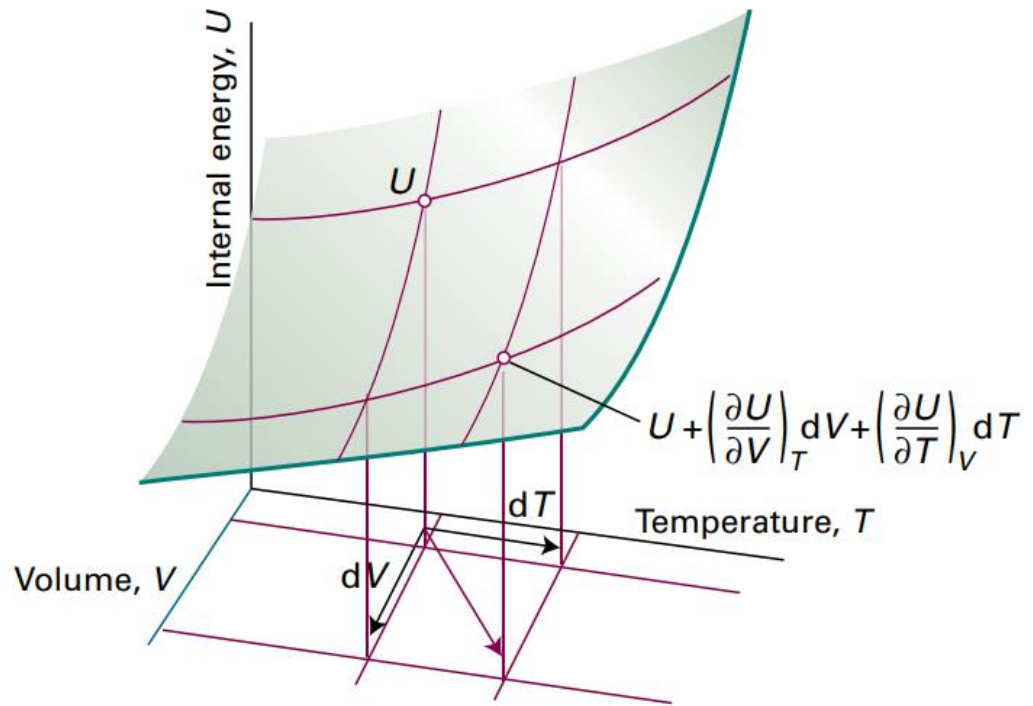
$$q = \int_{i, \text{path}}^f dq$$



inexact differential

A differential dX is called inexact if it represents the differential of a path function

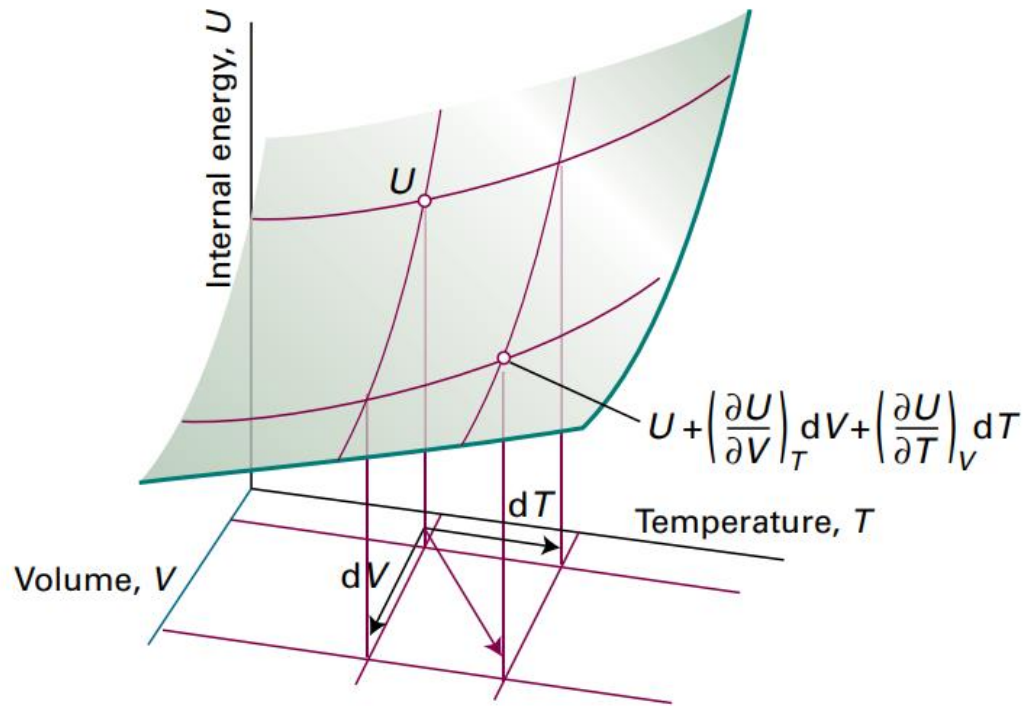
Changes in internal energy



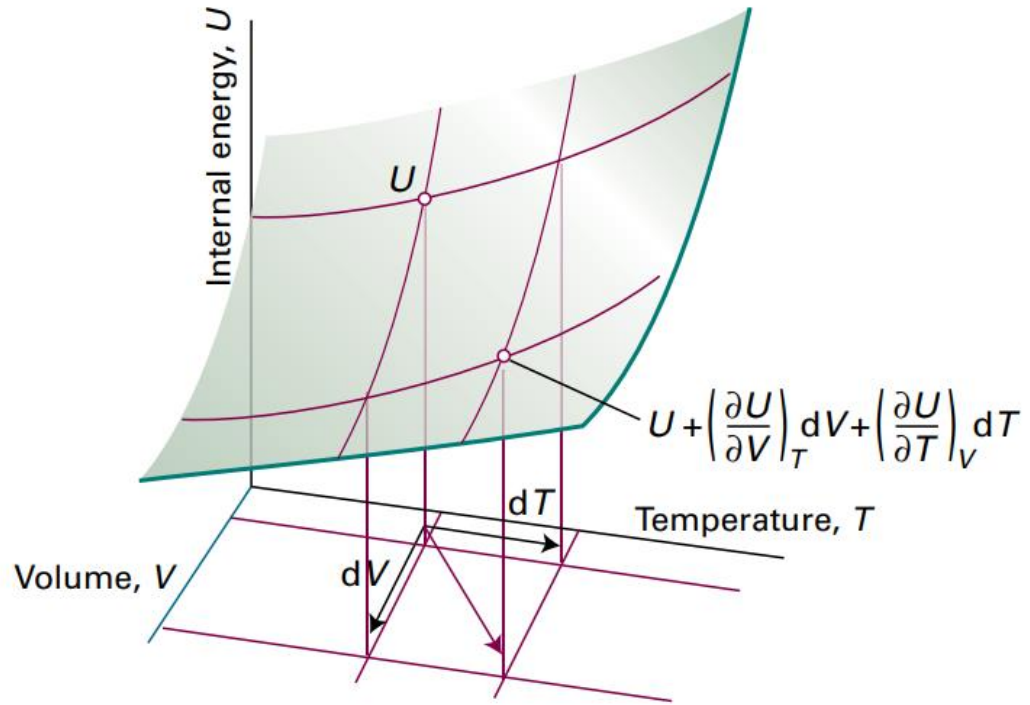
Changes in internal energy

If a function $U(V,T)$ depends on two independent variables v and T :

$$dU = \left(\frac{\partial U}{\partial V} \right)_T dV + \left(\frac{\partial U}{\partial T} \right)_V dT$$

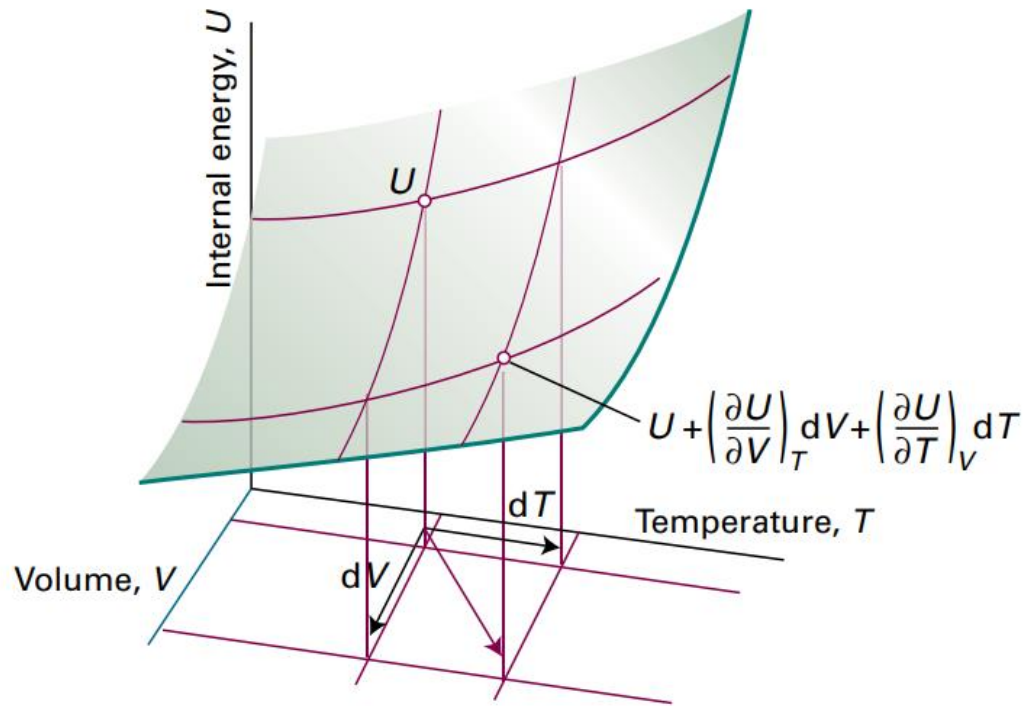


Changes in internal energy



$$dU = \left(\frac{\partial U}{\partial V}\right)_T dV + \overset{C_V}{\left(\frac{\partial U}{\partial T}\right)_V} dT$$

Changes in internal energy



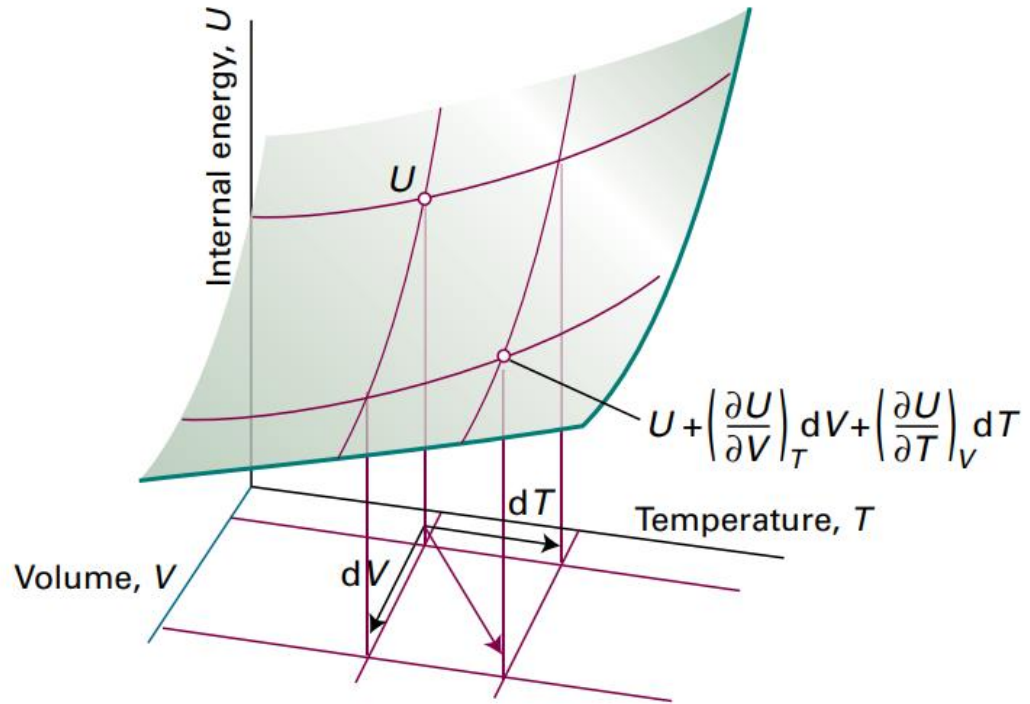
$$dU = \left(\frac{\partial U}{\partial V}\right)_T dV + \left(\frac{\partial U}{\partial T}\right)_V dT$$

↓

internal pressure

↑ C_V

Changes in internal energy

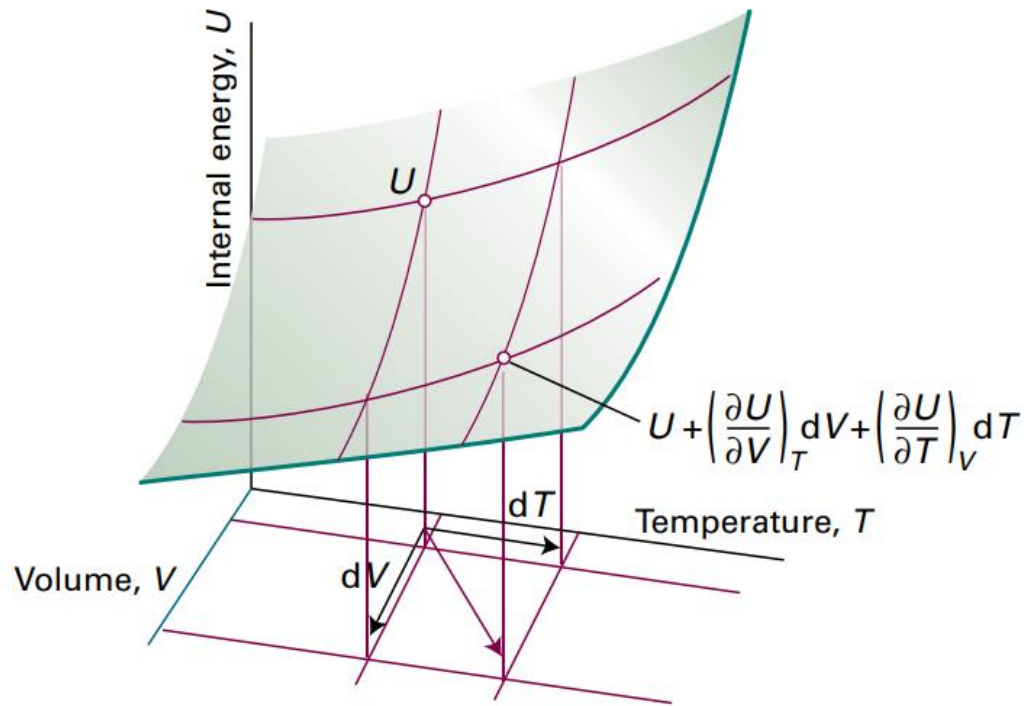


$$dU = \left(\frac{\partial U}{\partial V}\right)_T dV + \left(\frac{\partial U}{\partial T}\right)_V dT$$

internal pressure

$$dU = \pi_T dV + C_V dT$$

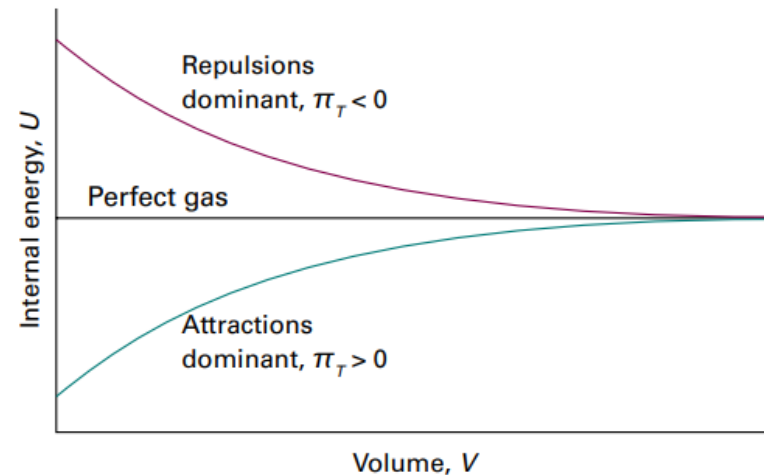
Changes in internal energy



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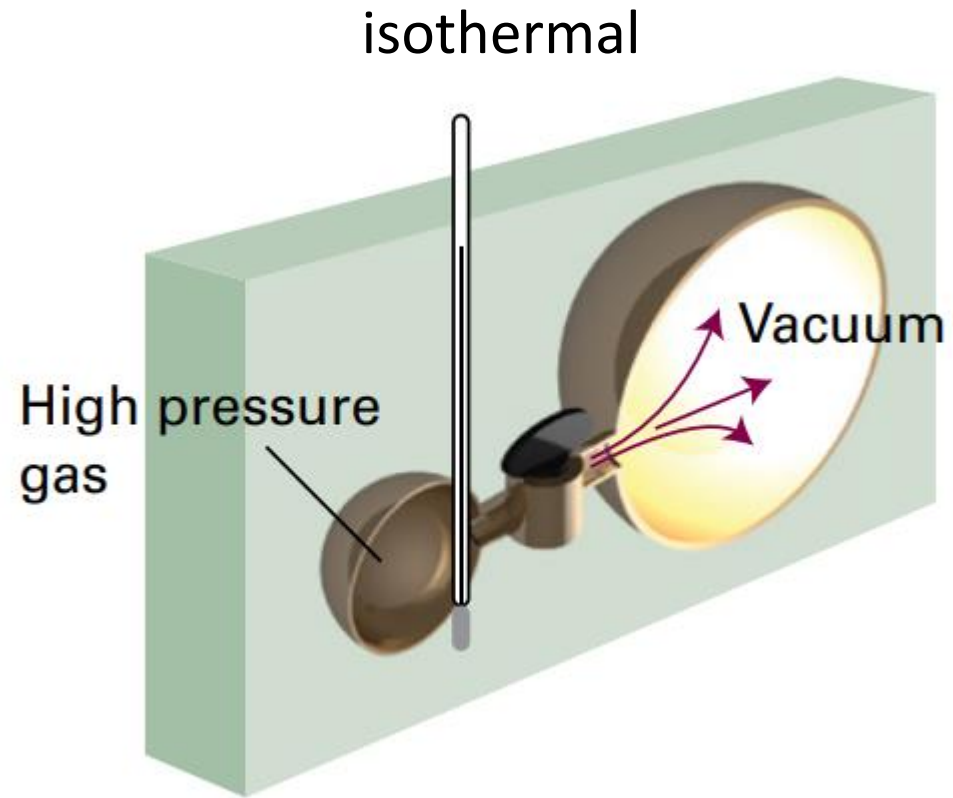
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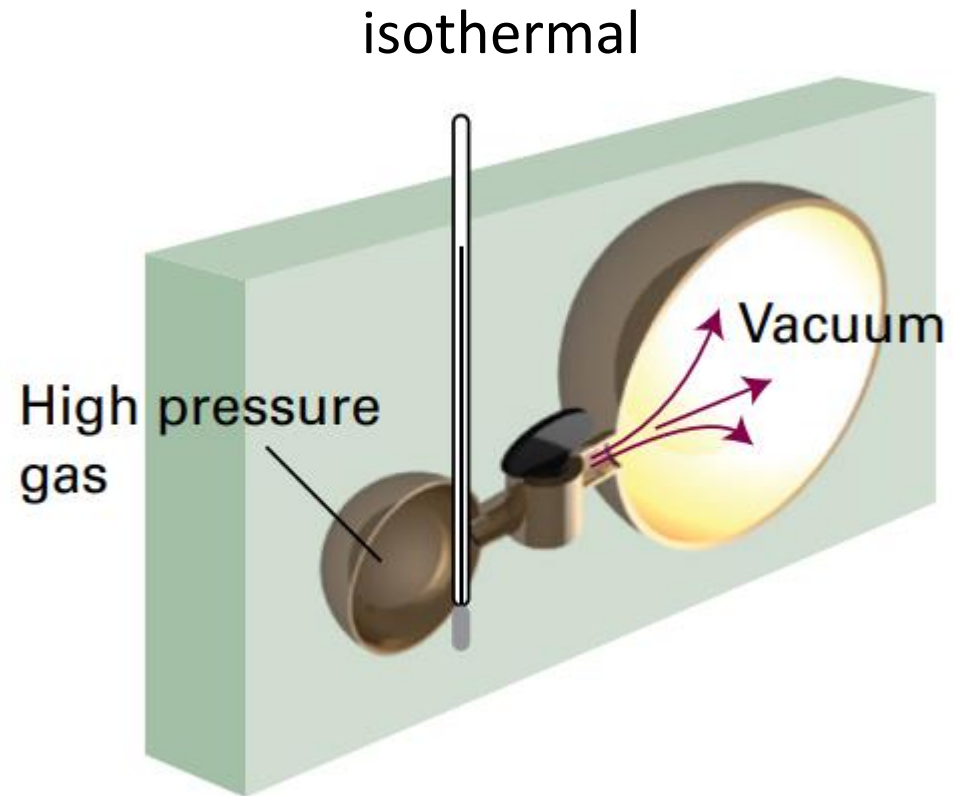
$\pi_T = 0$ for perfect gases

Measuring π_T



James Prescott Joule (1818 – 1889) was an English physicist

Measuring π_T



$$w = 0$$

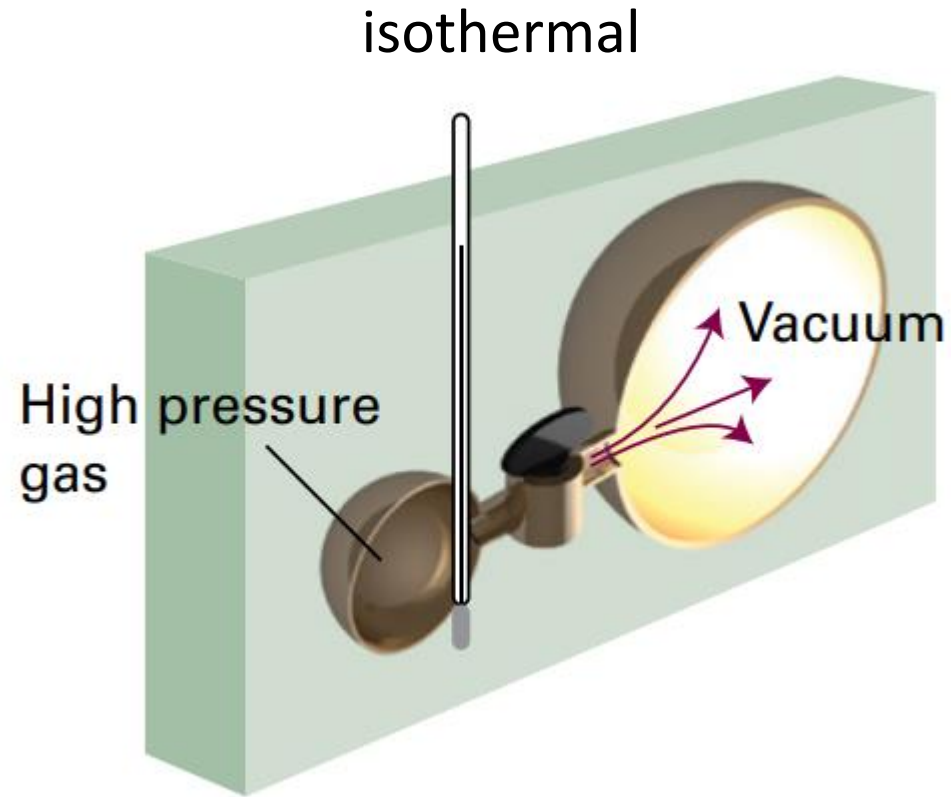
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$$\Delta U = 0$$



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Measuring π_T



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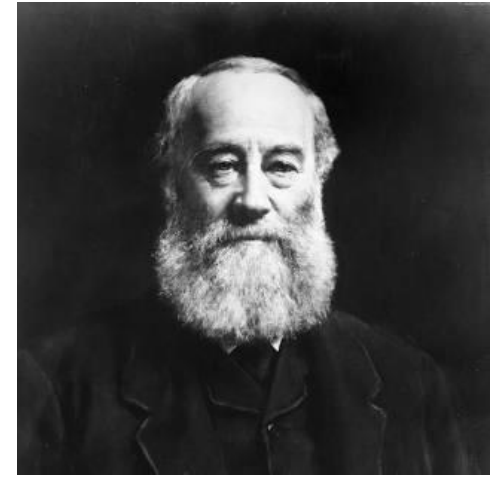
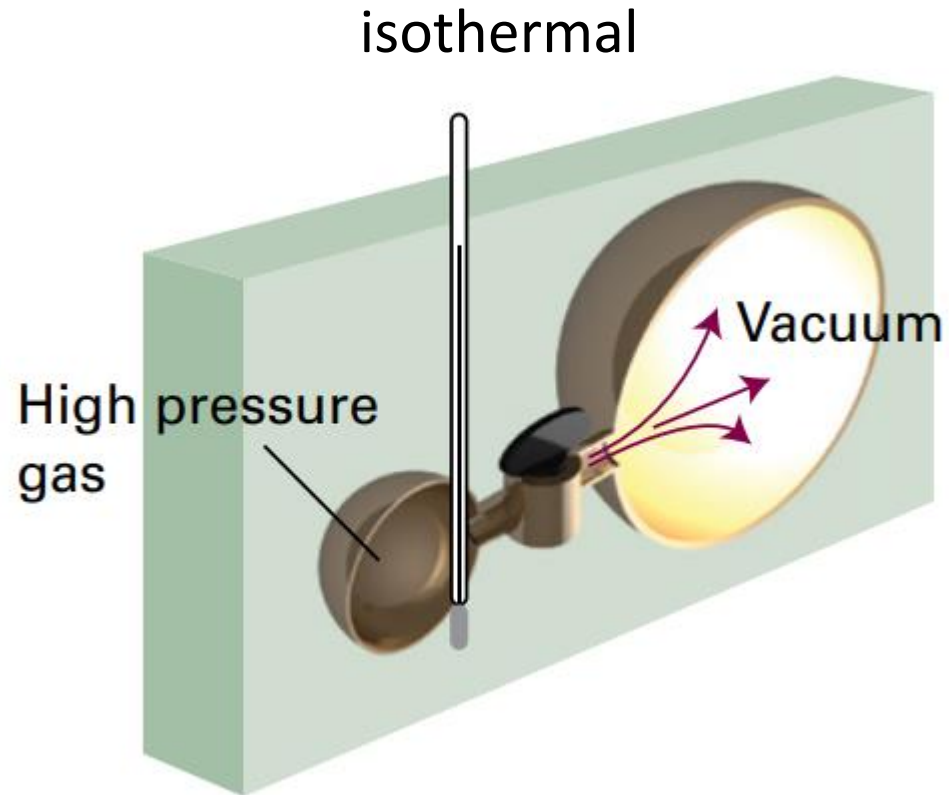
$$w = 0$$

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$$\text{No change in } U, \left(\frac{\partial U}{\partial V} \right)_T = \pi_T = 0$$

Measuring π_T



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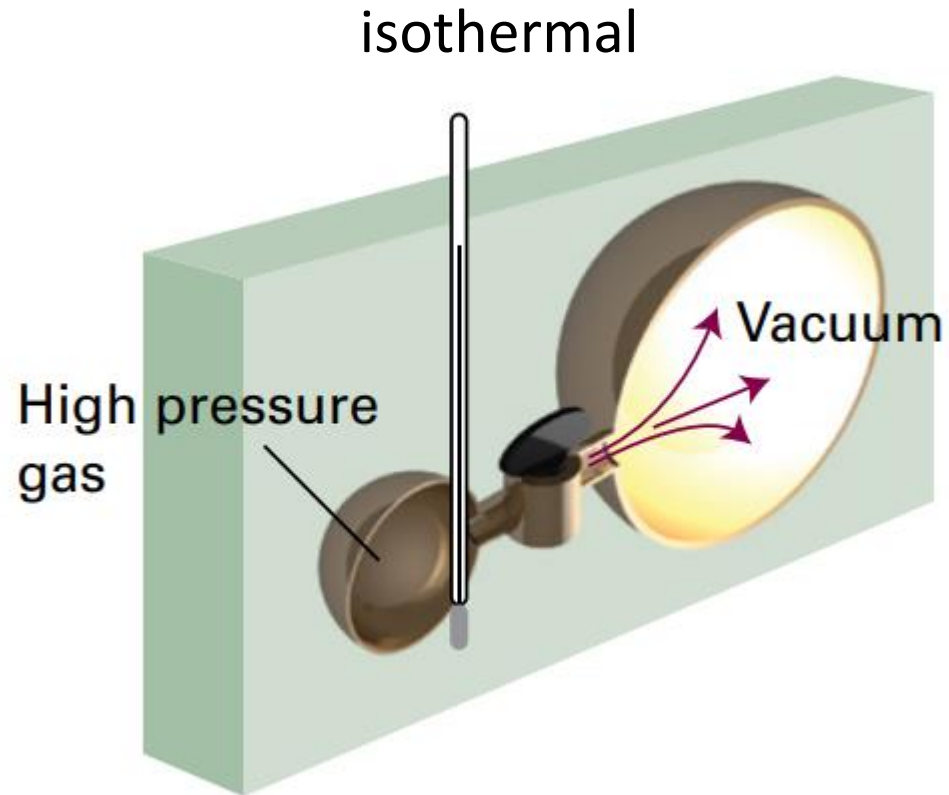
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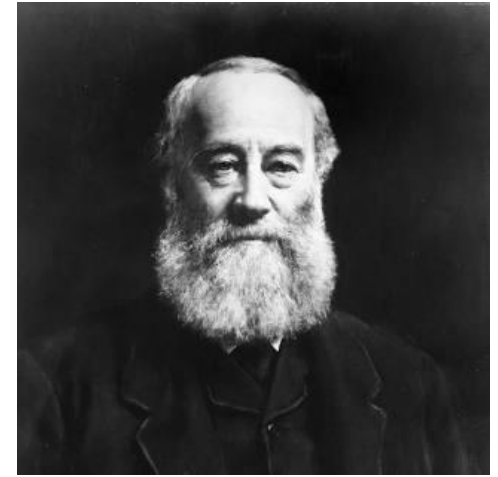
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Joule concluded that U does not change when a gas expands isothermally and therefore that $\pi_T = 0$.

Measuring π_T



Joule extracted a limited property of gas, a characteristic of a perfect gas, without noticing the small deviations in real gases.



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$$q = 0$$

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Expansion and isothermal coefficients

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P$$

expansion coefficient

Expansion and isothermal coefficients

	$\alpha/(10^{-4} \text{ K}^{-1})$	$\kappa_T/(10^{-6} \text{ bar}^{-1})$
<i>Liquids:</i>		
Benzene	12.4	90.9
Water	2.1	49.0
<i>Solids:</i>		
Diamond	0.030	0.185
Lead	0.861	2.18



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isothermal compressibility

Expansion and isothermal coefficients

For a perfect gas;

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p$$
$$= \frac{1}{V} \left(\frac{\partial (nRT/p)}{\partial T} \right)_p$$

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isothermal compressibility

Expansion and isothermal coefficients

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$$\begin{aligned}\alpha &= \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p \\ &= \frac{1}{V} \left(\frac{\partial (nRT/p)}{\partial T} \right)_p \\ &= \frac{1}{V} \times \frac{nR}{p} \\ &= \frac{nR}{nRT} = \frac{1}{T}\end{aligned}$$

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$$\kappa_T = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T$$

isothermal compressibility

For a perfect gas;

$$\kappa_T = 1/p$$